

The geometry of the vehicle is described by the wheelbase l and the distance of the center of gravity from the rear axle denoted by d . Moreover, the mass of the vehicle is denoted by m while its moment of inertia about its center of mass is denoted by J_G . The mass of the steering system (including the front wheel) is denoted by m_F while the moment of inertia about the center of the front wheel F is denoted by J_F . (The mass and the mass moment of inertia of the rear wheel are incorporated in m and J_G .)

We distinguish between the specific points of the wheels and tire-ground contact patches. In Figure 1, F and R denote the wheel centers for the front and rear tires, respectively. The leading points, where the center lines of the treads enter the contact regions are referred by Q and P, respectively. In order to eliminate the longitudinal dynamics we neglect the lateral deformation of the tires and consider that the driven wheels rotate with constant speed. Since we consider a front wheel drive car, we prescribe the wheel directional component of the velocity of the wheel center F to be a constant V . This poses a nonholonomic constraint on the dynamics of the system.

To derive the equations of motion we use the Appell-Gibbs formalism (Gantmacher, 1975),(De Sapio, 2017) which requires the introduction of the so-called pseudo velocities. Since the configuration is described by four generalized coordinates and there is one nonholonomic constraint, we need to choose three pseudo-velocities. Here we choose σ_1 to be the lateral velocity of the center of gravity G, σ_2 to be the yaw rate of the vehicle whereas σ_3 is the steering rate (the yaw rate of front axle relative to the car). The nonholonomic constraint and the definition of the pseudo-velocities can be summarized using the matrix equation

$$\begin{bmatrix} \cos(\psi + \delta) & \sin(\psi + \delta) & (l-d)\sin\delta & 0 \\ -\sin\psi & \cos\psi & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \dot{x} \\ \dot{y} \\ \dot{\psi} \\ \dot{\delta} \end{bmatrix} = \begin{bmatrix} V \\ \sigma_1 \\ \sigma_2 \\ \sigma_3 \end{bmatrix}. \quad (1)$$

Solving this equation the time derivatives of the generalised coordinates can be expressed as

$$\begin{aligned} \dot{x} &= V \frac{\cos\psi}{\cos\delta} - \sigma_1 \frac{\sin(\psi + \delta)}{\cos\delta} - \sigma_2(l-d)\cos\psi \tan\delta, \\ \dot{y} &= V \frac{\sin\psi}{\cos\delta} + \sigma_1 \frac{\cos(\psi + \delta)}{\cos\delta} - \sigma_2(l-d)\sin\psi \tan\delta, \\ \dot{\psi} &= \sigma_2, \\ \dot{\delta} &= \sigma_3. \end{aligned} \quad (2)$$

These kinematic equations give part of the governing equations, while the other part is comprised of the Appell-Gibbs equations:

$$\begin{bmatrix} m_{11} & m_{12} & 0 \\ m_{21} & m_{22} & J_F \\ 0 & J_F & J_F \end{bmatrix} \begin{bmatrix} \dot{\sigma}_1 \\ \dot{\sigma}_2 \\ \dot{\sigma}_3 \end{bmatrix} = \begin{bmatrix} f_1 \\ f_2 \\ f_3 \end{bmatrix}, \quad (3)$$

where the elements of the generalised mass matrix are given by

$$\begin{aligned} m_{11} &= \frac{m_F + m}{\cos^2\delta}, \\ m_{12} = m_{21} &= \frac{m_F + m \sin^2\delta}{\cos^2\delta}(l-d), \\ m_{22} &= J_F + J_G + \frac{m_F + m \sin^2\delta}{\cos^2\delta}(l-d)^2, \end{aligned} \quad (4)$$

while the right hand side can be expressed as

$$\begin{aligned} f_1 &= \frac{F_F}{\cos\delta} + F_R \\ &+ \frac{1}{\cos\delta} \left(-(m_F + m)V + m\sigma_2(l-d)\sin\delta \right) \sigma_2 \\ &+ \frac{(m_F + m)\sin\delta}{\cos^3\delta} \left(V\sin\delta - \sigma_1 - (l-d)\sigma_2 \right) \sigma_3, \\ f_2 &= M_F + M_R + M_S + \frac{(l-d)F_F}{\cos\delta} - dF_R \\ &- \frac{l-d}{\cos\delta} \left(m_F V + m\sigma_1 \sin\delta \right) \sigma_2 \\ &+ \frac{(m_F + m)(l-d)\sin\delta}{\cos^3\delta} \left(V\sin\delta - \sigma_1 - (l-d)\sigma_2 \right) \sigma_3, \\ f_3 &= M_F + M_S. \end{aligned} \quad (5)$$

Here M_S denotes the steering torque while F_F , F_R and M_F , M_R are the lateral tire forces and aligning torques calculated from the tire model. In particular, for the brush tire model these can be expressed by the formulae

$$F(\alpha) = \begin{cases} \phi_3 \tan^3 \alpha + \phi_2 \tan^2 \alpha \operatorname{sgn} \alpha + \phi_1 \tan \alpha, & 0 \leq |\alpha| < \alpha_{\text{crit}}, \\ \mu F_z \operatorname{sgn} \alpha, & \alpha_{\text{crit}} < |\alpha|, \end{cases} \quad (6)$$

and

$$M(\alpha) = \begin{cases} \mu_4 \tan^4 \alpha \operatorname{sgn} \alpha + \mu_3 \tan^3 \alpha + \mu_2 \tan^2 \alpha \operatorname{sgn} \alpha + \mu_1 \tan \alpha, & 0 \leq |\alpha| < \alpha_{\text{crit}}, \\ 0, & \alpha_{\text{crit}} < |\alpha|, \end{cases} \quad (7)$$

as function of the slide-slip angle α .

The slide slip angles can be calculated from vehicle kinematics as

$$\begin{aligned} \tan \alpha_F &= \frac{\sigma_1 + (l-d)\sigma_2 + a(\sigma_2 + \sigma_3)}{V \cos\delta} - \tan\delta, \\ \tan \alpha_R &= \frac{(\sigma_1 - (d-a)\sigma_2) \cos\delta}{V - (\sigma_1 + (l-d)\sigma_2) \sin\delta}, \end{aligned} \quad (8)$$

for the front and the rear wheels, respectively. In the formulae above μ denotes the friction coefficient, F_z is the vertical load on the tire while α_{crit} corresponds to the critical side-slip angle at which the whole contact patch starts to slide. Moreover, the constants ϕ_i and μ_i are given in Appendix A and a denotes the half-length of the tire-ground contact patch (i.e., the distance between F and Q and the distance between R and P).

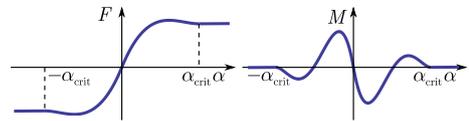


Fig. 2. Qualitative lateral tire force (left panel) and aligning torque (right panel) characteristics.

Formulae (6) and (7) are visualized in Fig. 2. While these characteristics are piecewise-smooth continuous functions, for linear stability analysis of the rectilinear motion it is sufficient to use the continuous linear parts for small α , that is,

$$F(\alpha) \approx \phi_1 \alpha = 2a^2 k \alpha, \quad (9)$$

and

$$M(\alpha) \approx \mu_1 \alpha = -\frac{2}{3} a^3 k \alpha, \quad (10)$$

where a is the half-length of the tire-ground contact patch while k is the lateral stiffness of the tire distributed along the contact length; see the definition of ϕ_1 and μ_1 in (A.1) and (A.4) in Appendix A. The constant $2a^2k$ in (10) is often referred as the cornering stiffness.

Here we construct a hierarchical controller in order to regulate the steering torque M_S such that the vehicle follows a straight path. Without loss of generality we make the vehicle to run along the x axis, i.e., $y(t) \equiv 0$ and $\psi(t) \equiv 0$. At the higher-level we calculate a desired steering angle from the vehicle position and orientation as

$$\delta_{\text{des}}(t) = -k_\psi \sin(\psi(t - \tau_1)) - k_y y(t - \tau_1), \quad (11)$$

where k_ψ and k_y are proportional gains for the lateral position and the orientation angle while a time delay τ_1 represent the time needed for sensing, computation, and actuation. Then the steering torque is determined by the PID controller

$$M_S(t) = k_p(\delta_{\text{des}}(t - \tau_2) - \delta(t - \tau_2)) + k_d(\dot{\delta}_{\text{des}}(t - \tau_2) - \dot{\delta}(t - \tau_2)) + k_i z(t - \tau_2), \quad (12)$$

where τ_2 represent the time delay in the control loop and the integral term is considered through the variable z defined by

$$\dot{z}(t) = \delta_{\text{des}}(t) - \delta(t). \quad (13)$$

In our study we consider the gains of the lower-level controller in the form of

$$k_p = pk_{p0}, \quad (14)$$

$$k_d = pk_{d0}, \quad (15)$$

$$k_i = pk_{i0}. \quad (16)$$

The benefit of this formulation is that parameter p can be used to represent the ‘strength’ of the lower-level controller for a fixed set of the gains k_{p0} , k_{d0} and k_{i0} . Note that $p = 0$ corresponds to the case of no control on the front axle, while $p \rightarrow \infty$ with $\tau_2 = 0$ corresponds to the ideal case when the steering angle can be directly set.

3. LINEAR STABILITY ANALYSIS

The closed-loop dynamics of the vehicle is described by the equations (2,3,11,12,13) that can be written into the compact form

$$\dot{\mathbf{x}}(t) = \boldsymbol{\phi}(\mathbf{x}(t), \mathbf{x}(t - \tau_1), \mathbf{x}(t - \tau_1 - \tau_2)), \quad (17)$$

where

$$\mathbf{x} = [x \ y \ \psi \ \delta \ \sigma_1 \ \sigma_2 \ \sigma_3 \ z]^T, \quad (18)$$

contains the state variables. The rectilinear motion can be expressed as

$$\mathbf{x}^*(t) = [Vt \ 0 \ 0 \ 0 \ 0 \ 0 \ 0 \ 0]^T. \quad (19)$$

Let us introduce the perturbations $\tilde{\mathbf{x}}(t) = \mathbf{x}(t) - \mathbf{x}^*(t)$ and linearize (17) about the the rectilinear motion to obtain

$$\dot{\tilde{\mathbf{x}}}(t) = \mathbf{A}_0 \tilde{\mathbf{x}}(t) + \mathbf{A}_{\tau_1} \tilde{\mathbf{x}}(t - \tau_1) + \mathbf{A}_{\tau_2} \tilde{\mathbf{x}}(t - \tau_1 - \tau_2), \quad (20)$$

where the matrixes \mathbf{A}_0 , \mathbf{A}_{τ_1} , and \mathbf{A}_{τ_2} are defined in Appendix B.

Using the trial solution $\mathbf{x}(t) = \mathbf{c}e^{\lambda t}$ we can obtain the characteristic equation as

$$D(\lambda) = \det \left(\mathbf{A}_0 + \mathbf{A}_{\tau_1} e^{-\lambda \tau_1} + \mathbf{A}_{\tau_2} e^{-\lambda(\tau_1 + \tau_2)} \right) = 0, \quad (21)$$

which has infinitely many solutions for the characteristic roots λ . Notice that $\lambda = 0$ is always a solution corresponding to the ‘neutral’ direction x . After eliminating this

trivial eigenvalue by defining $\tilde{D}(\lambda) = D(\lambda)/\lambda$, $\tilde{D}(0) = 0$ provides the stability boundaries for non-oscillatory stability loss whereas $\tilde{D}(i\omega) = 0$ gives the stability boundaries for oscillatory stability loss. These bound the stable domain in the parameter space where all eigenvalues have negative real part. When crossing a $\tilde{D}(0) = 0$ boundary a real characteristic root moves to the right half complex plane, while when crossing the $\tilde{D}(i\omega) = 0$ boundary a pair of complex conjugate roots crosses the imaginary axis from left to right (Stépán, 1989). In this case oscillations arise with angular frequency close to ω .

In order to calculate the stability boundaries numerically we use semi-discretisation (Insperger and Stépán, 2013). In particular, we utilize Chebyshev polynomials and the Chebyshev differentiation matrix to approximate the infinite dimensional spectrum of (20); see (Trefethen, 2000). Using these techniques we draw stability charts in the plane of the control gains for different values of the delays.

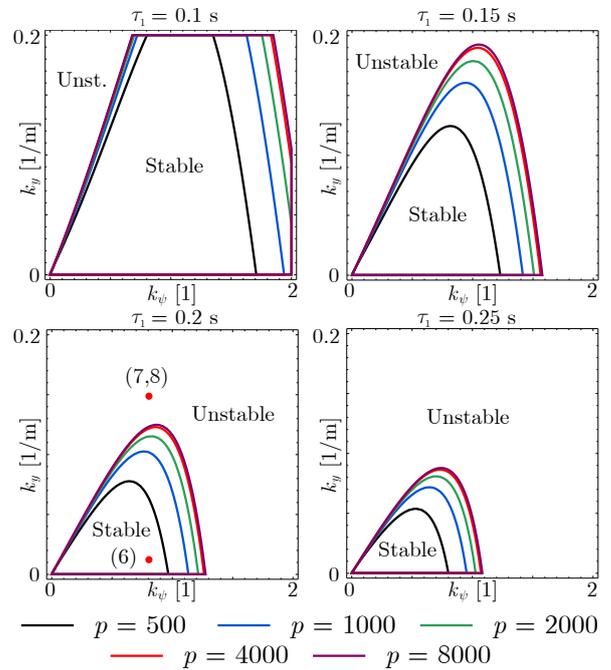


Fig. 3. Stability charts for different values of higher-level delay τ_1 in the plane of the higher-level control gains k_ψ and k_y with lower-level delay $\tau_2 = 0.0001$ s. The boundaries corresponding to different values of control strength p are shown in different colours. In the top left panel parts of stable regions are outside the window.

In Fig. 3 the stability boundaries are shown for different higher-level delays τ_1 in the plane of the higher-level control gains (k_ψ, k_y) corresponding to the lateral position and orientation of the vehicle. The different curves in each panel correspond to different values of parameter p while the lower-level delay is set to $\tau_2 = 0.0001$ s which is small enough not to endanger stability. Then the lower-level control gains are set to $k_{p0} = 8$ Nm, $k_{d0} = 0.1$ Nms, $k_{i0} = 0.5$ Nm/s while the vehicle and tire parameters are listed in Table 1 and they are kept the same in the rest of the paper. It can be seen, that the size of the stable parameter domain strongly depends on the time delay in the higher-level

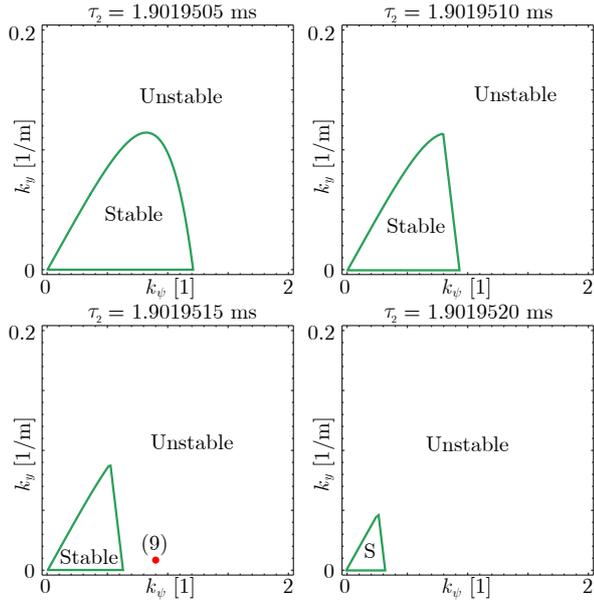


Fig. 4. Stability charts for different values of lower-level delay τ_2 in the plane of the higher-level control gains k_ψ and k_y with higher-level delay of $\tau_1 = 0.2$ s and control strength of $p = 2000$ are considered.

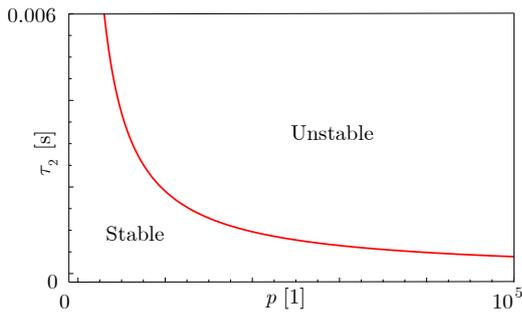


Fig. 5. Stability chart in the plane of the control 'strength' p and the time delay τ_2 of the lower-level controller. The higher level gains are set to $k_\psi = 0.5$ and $k_y = 0.01$ 1/m whereas the higher-level delay of $\tau_1 = 0.2$ s is considered.

control loop, i.e., the larger the delay is, the smaller the stable parameter domain is. The larger values of parameter p result in larger stable domains. However, this growth has an upper limit corresponding to $p \rightarrow \infty$, i.e., the 'ideal' lower-level controller. (Practically the value $p = 8000$ can be regarded as an almost 'ideal' lower-level controller).

Table 1. Vehicle and tire parameters

Parameter name	Value	Unit
l wheelbase	2.57	m
d center of gravity position	1.54	m
m vehicle mass	1100	kg
J_G vehicle mass moment of inertia	1343	kgm ²
m_F front axle mass	10	kg
J_F front axle mass moment of inertia	0.25	kgm ²
V longitudinal velocity	15	m/s
a tire-ground contact half-length	0.1	m
k lateral tire stiffness	2×10^6	N/m

If the delay τ_2 of the lower-level controller is non-zero but sufficiently small it has a negligible effect on the stability boundaries. However, within a very narrow interval of this delay the whole stable parameter region disappears, as shown in Fig. 4 where the lower-level delay is varied only by a fraction of a millisecond. This means that the critical value of the lower-level delay is only very weakly dependent to the parameters of the higher-level controller. Thus, the two controllers can be designed and investigated separately. In the meantime, it can be established that the critical lower-level delay strongly depends on the control strength p as shown in Fig. 5.

The investigation of the imaginary parts of the critical eigenvalue reveals that the frequency corresponding to the arising vibrations are considerably different in the two control levels. If the higher-level control gains are chosen from the unstable parameter domain vibrations with a frequency up to ~ 3 -5 Hz can be observed. On the contrary due to instabilities of the lower-level controller one can observe high-frequency vibrations (~ 100 -150 Hz).

4. SIMULATIONS

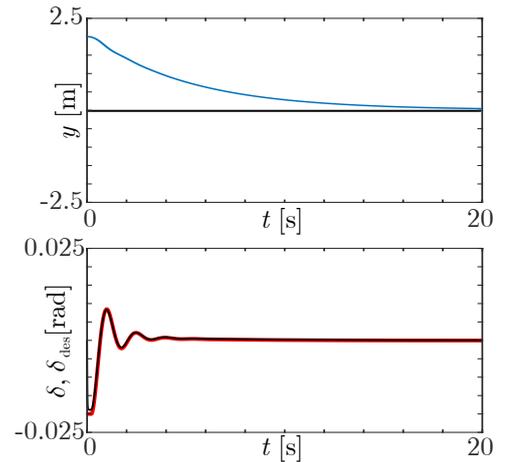


Fig. 6. Time profiles for the centre of gravity's lateral position (upper panel) and the desired (thick red) and actual (thin black) steering angles (lower panel) for control parameters $k_\psi = 0.8$, $k_y = 0.01$ 1/m, $p = 2000$ and time-delays $\tau_1 = 0.1$ s, $\tau_2 = 0.0001$ s.

Numerical simulations have been carried out for the non-linear system to demonstrate the behavior for different scenarios. In Figure 6 the lateral position of the vehicle centre of gravity and the desired and actual steering angles are presented for a stable set of control parameters shown in the bottom left panel of Fig. 3. In Figures 7 and 8 we demonstrate the developing oscillations due to 'badly' chosen higher-level control gains. The two different initial conditions are applied to show that the system converges to a stable limit cycle. In Figure 9 we present the system behavior in the case when the delay in the lower-level controller is larger than the critical one, see the point marked on the bottom left panel in Fig. 4.

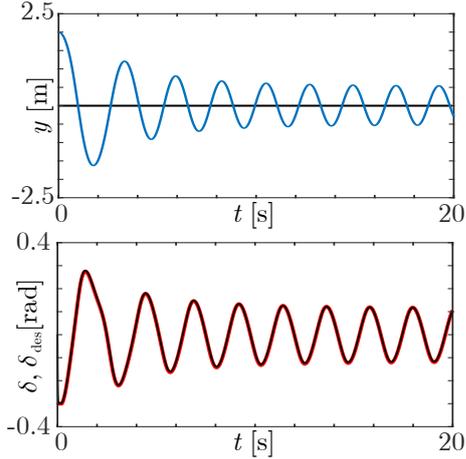


Fig. 7. Time profiles for the centre of gravity's lateral position (upper panel) and the desired (thick red) and actual (thin black) steering angles (lower panel) for control parameters $k_\psi = 0.8$, $k_y = 0.15$ 1/m, $p = 2000$ and time-delays $\tau_1 = 0.1$ s, $\tau_2 = 0.0001$ s.

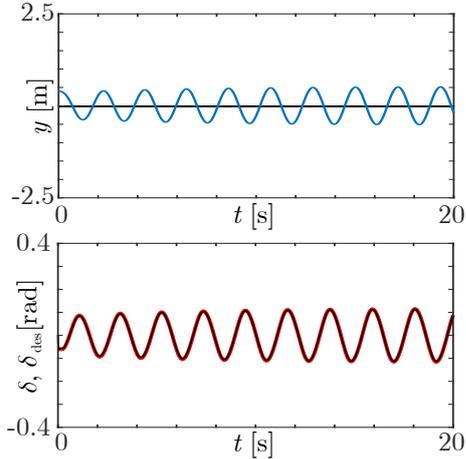


Fig. 8. Time profiles for the centre of gravity's lateral position (upper panel) and the desired (thick red) and actual (thin black) steering angles (lower panel) for control parameters $k_\psi = 0.8$, $k_y = 0.15$ 1/m, $p = 2000$ and time-delays $\tau_1 = 0.1$ s, $\tau_2 = 0.0001$ s.

5. CONCLUSIONS

In the paper a hierarchical lateral controller of an autonomous vehicle was investigated considering different time delays in the two control levels. The main result of the analysis is, that although the higher- and the lower-level controllers are designed to operate together, their influence on each other is weak. In practice the control gains of the higher-level controller may be chosen by assuming that the desired steering angle can be set accurately. Similarly, if the gains are large enough the stability of the lower-level controller is almost independent from the way the higher-level algorithm is designed. As we demonstrated, the time delay τ_2 in the lower-level controller is a critical parameter from the point of view of the stability. While the higher-level controller can be stable even for relatively large delays, at the lower-level even a delay of a couple

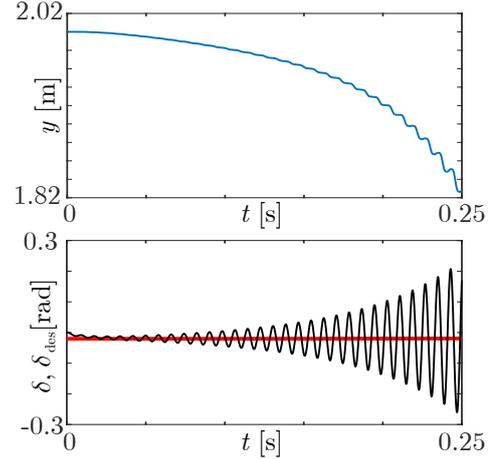


Fig. 9. Time profiles for the centre of gravity's lateral position (upper panel) and the desired (thick red) and actual (thin black) steering angles (lower panel) for control parameters $k_\psi = 0.8$, $k_y = 0.01$ 1/m, $p = 2000$ and time-delays $\tau_1 = 0.1$ s, $\tau_2 = 0.0019019505$ s.

of milliseconds can result unstable behavior. Therefore it is required that such algorithms run with a sampling frequency around 1000 Hz.

Further analysis could be carried out considering a discrete-time system to emulate the digital controllers more accurately. The nonlinear analysis of the system can be subject of further studies, too. On one hand, considering the geometrical nonlinearities as well as the nonlinear piecewise-smooth tire characteristics might reveal 'bistable' parameter domains where a stable equilibrium and a stable periodic solution coexist. This behavior considered to be dangerous is dangerous from engineering point of view as the domain of attraction of the stable equilibrium can be limited. On the other hand, one may need to consider possible saturations of the actuators which can also lead to complex dynamical phenomena.

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Appendix A. COEFFICIENTS IN THE TIRE-FORCE CHARACTERISTICS

The coefficients in the tire force and self aligning torque characteristics in (6) and (7) can be expressed as

$$\phi_1 = 2a^2k, \quad (\text{A.1})$$

$$\phi_2 = -\frac{8a^4k^2}{3F_z\mu_0} + \frac{4a^4k^2\mu}{3F_z\mu_0^2}, \quad (\text{A.2})$$

$$\phi_3 = -\frac{8a^6k^3}{9F_z^2\mu_0^2} - \frac{16a^6k^3\mu}{27F_z^2\mu_0^3} \quad (\text{A.3})$$

and

$$\mu_1 = -\frac{2}{3}a^3k, \quad (\text{A.4})$$

$$\mu_2 = +\frac{8a^5k^2}{3F_z\mu_0} - \frac{4a^5k^2\mu}{3F_z\mu_0^2}, \quad (\text{A.5})$$

$$\mu_3 = -\frac{8a^7k^3}{9F_z^2\mu_0^2} + \frac{16a^7k^3\mu}{9F_z^2\mu_0^3}, \quad (\text{A.6})$$

$$\mu_4 = +\frac{64a^9k^4}{81F_z^3\mu_0^3} - \frac{16a^9k^4\mu}{9F_z^3\mu_0^4}, \quad (\text{A.7})$$

where a is the contact patch half-length, k is the distributed lateral stiffness of the tires, F_z is the vertical load on the axles, μ and μ_0 are the friction coefficients between the tires and the road for sliding and rolling, respectively.

Appendix B. COEFFICIENT MATRICES OF THE LINEARIZED SYSTEM

The coefficient matrices in (21) can be expressed as

$$\mathbf{A}_0 = \begin{bmatrix} a_{11} & a_{12} & a_{13} & 0 & 0 & 0 & a_{17} & 0 \\ a_{21} & a_{22} & a_{23} & 0 & 0 & 0 & a_{27} & 0 \\ a_{31} & a_{32} & a_{33} & 0 & 0 & 0 & a_{37} & \frac{k_i}{J_F} \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & V & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}, \quad (\text{B.1})$$

where

$$a_{11} = -\frac{2a^2k(6J_G + (a-3l)(d-l)m_F)}{3V\Delta}, \quad (\text{B.2})$$

$$a_{12} = -\left((2ka^4 - 2k(d+3l)a^3 + 6kdl a^2 + 3(d-l)mV^2)(d-l)m_F + 3(4ka^3 + 2k(l-2d)a^2 + (m+m_F)V^2)J_G\right)/(3V\Delta), \quad (\text{B.3})$$

$$a_{13} = -\frac{2a^3kJ_G}{V\Delta}, \quad (\text{B.4})$$

$$a_{17} = \frac{2a^2kJ_G}{\Delta}, \quad (\text{B.5})$$

$$a_{21} = -\frac{2a^2k((a-6d+3l)m + (a-3l)m_F)}{3V\Delta}, \quad (\text{B.6})$$

$$a_{22} = -2a^2k((a^2 - 7da + 3la + 6d^2 + 3l^2 - 6dl)m + (a-d)(a-3l)m_F)/(3V\Delta), \quad (\text{B.7})$$

$$a_{23} = \frac{2a^3k(d-l)m}{V\Delta}, \quad (\text{B.8})$$

$$a_{27} = \frac{2a^2k(l-d)m}{\Delta}, \quad (\text{B.9})$$

$$a_{31} = \frac{2a^2k(((a-3l)m_F + (a-6d+3l)m)J_F - a\Delta)}{3V\Delta J_F} \quad (\text{B.10})$$

$$a_{32} = 2a^2k\left(\left((a^2 - 7da + 3la + 6d^2 + 3l^2 - 6dl)m + ((a-d)(a-3l)m_F)J_F - a(a-d+l)\Delta\right)/(3V\Delta J_F)\right), \quad (\text{B.11})$$

$$a_{33} = -\frac{2a^3k(3(d-l)mJ_F + a\Delta)}{3V\Delta J_F}, \quad (\text{B.12})$$

$$a_{37} = \frac{2a^2k(3(d-l)mJ_F + a\Delta)}{3V\Delta J_F}, \quad (\text{B.13})$$

and

$$\Delta = (d-l)^2m_Fm + (m_F + m)J_G. \quad (\text{B.14})$$

Moreover, we have

$$\mathbf{A}_{\tau_1} = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -\frac{k_d}{J_F} & 0 & 0 & 0 & -\frac{k_p}{J_F} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 \end{bmatrix}, \quad (\text{B.15})$$

$\mathbf{A}_{\tau_2} =$

$$\begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ -\frac{k_d k_y}{J_F} & -\frac{k_d k_\psi}{J_F} & 0 & 0 & -\frac{k_p k_y}{J_F} & -\frac{V k_d k_y + k_p k_\psi}{J_F} & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -k_y & -k_\psi & 0 & 0 \end{bmatrix}. \quad (\text{B.16})$$